

Directed polymer in γ -stable Random Environments

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Chemical background

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Polymers abound in nature because of the multivalency of atoms like carbon, silicon, oxygen, nitrogen, sulfur and phosphorus, which are capable of forming long concatenated structures.

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- linear, branched.

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- can wind around itself to form knots
- can be extended due to repulsive forces between the monomers as a result of excluded-volume
- or can collapse to a ball due to attractive forces.

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- Free energy of the polymer in this limit,
- presence of phase transitions as a function of underlying model parameters

Random Polymer

\mathbf{P}_x : Probability measure on $(\Omega, \mathcal{F}) := \left((\mathbb{Z}^d)^{\mathbb{N}}, \mathcal{P}(\mathbb{Z}^d)^{\otimes \mathbb{N}} \right)$ of sequences $S := (S_n)_{n \geq 0}$ such that:

$$\begin{aligned} S_0 &= x, \\ \{S_n - S_{n-1}\}_{n \geq 1} &\text{ is an IID sequence, and} \\ \mathbf{P}_x[S_1 = x + e_j] &= \mathbf{P}_x[S_1 = x - e_j] = \frac{1}{2d}, \end{aligned} \tag{1}$$

Random Environment

IID random variables $\eta := \{\eta_{n,z} : n \in \mathbb{N}, z \in \mathbb{Z}^d\}$, called *the environment*, defined on a probability space $(\Lambda, \mathcal{F}, \mathbb{P})$, that satisfies:

$$\begin{aligned} \mathbb{E}[\eta_{0,0}] &= 0 \text{ and} \\ \mathbb{E}[\exp(\beta\eta_{0,0})] &< \infty, \text{ for all } \beta \in \mathbb{R}. \end{aligned} \tag{2}$$

Polymer Measure

Given $\beta > 0$, $N \in \mathbb{N}$ and a fixed realization of the environment η , we define the measure $\mathbf{P}_N^{\beta, \eta}$ on the space Ω , called the *polymer measure*, by its Radon-Nikodym derivative with respect to \mathbf{P}_0 :

$$\frac{d\mathbf{P}_N^{\beta, \eta}}{d\mathbf{P}_0}(S) = \frac{1}{Z_N^{\beta, \eta}} \exp \left(\beta \sum_{n=1}^N \eta_{n, S_n} \right), \quad (3)$$

where $Z_N^{\beta, \eta}$ is the positive normalization factor that makes $\mathbf{P}_N^{\beta, \eta}$ a probability measure.

Bolthausen '87

$$W_N^{\beta,\eta} := \frac{Z_N^{\beta,\eta}}{\mathbb{E} \left[Z_N^{\beta,\eta} \right]}, \quad (4)$$

$$W_\infty^{\beta,\eta} := \lim_{N \rightarrow \infty} W_N^{\beta,\eta}, \quad (5)$$

exists \mathbb{P} -a.s. and is a non-negative random variable.

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we have

- *weak disorder* if $W_\infty^\beta > 0$ \mathbb{P} -a.s. and
- *strong disorder* if $W_\infty^\beta = 0$ \mathbb{P} -a.s..

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Comets, Yoshida 2005

Assuming $d \geq 3$ and weak disorder, the measures $\mathbf{P}_N^{\beta, \eta}$, after rescaling, converge in law to the Brownian Motion, for almost all realizations of the environment.

Strong disorder

The polymer is largely influenced by the disorder and is attracted to sites with favorable environment (localize phase).

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Comtets, Shiga, Yoshida 2003

$$\left\{ W_{\infty}^{\beta, \eta} = 0 \right\} = \left\{ \sum_{n \geq 1} \left(\mathbf{P}_{n-1}^{\beta, \eta} \right)^{\otimes 2} [S_n = S'_n] = \infty \right\} \mathbb{P}\text{-a.s.}, \quad (6)$$

where S and S' are two independent polymers with distribution $\mathbf{P}_{n-1}^{\beta, \eta}$. Moreover, if $\mathbb{P}[W_{\infty}^{\beta, \eta} = 0] = 1$, then there exists some constants $c_1, c_2 \in (0, \infty)$ such that,

$$-c_1 \log W_N^{\beta, \eta} \leq \sum_{n \geq 1} \left(\mathbf{P}_{n-1}^{\beta, \eta} \right)^{\otimes 2} [S_n = S'_n] \leq -c_2 \log W_N^{\beta, \eta} \quad (7)$$

Comets, Yoshida 2005

There exists a critical value $\beta_c = \beta_c(d) \in [0, \infty]$ with

$$\beta_c = 0 \text{ for } d = 1, 2 \text{ and} \quad (8)$$

$$\beta_c > 0 \text{ for } d \geq 3, \quad (9)$$

such that there is weak disorder for $\beta \in [0, \beta_c)$ and strong disorder for $\beta > \beta_c$.

Example: Showing weak disorder for $d \geq 3$ and small β .

$$\mathbb{E} \left[\left(W_N^{\beta, \eta} \right)^2 \right] = \mathbb{E} \mathbf{E}^{\otimes 2} \left[\exp \left(\sum_{i \geq N} \beta \eta_{i, s_i} - \frac{\beta^2}{2} + \beta \eta_{i, s'_i} - \frac{\beta^2}{2} \right) \right] \quad (10)$$

$$\leq \mathbf{E}^{\otimes 2} \left[\exp \left(\sum_i \beta^2 \mathbf{1}_{\{s_i = s'_i\}} \right) \right]. \quad (11)$$

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$$\leq \mathbf{E}^{\otimes 2} \left[\exp \left(\sum_i \beta^2 \mathbf{1}_{\{S_i = S'_i\}} \right) \right]. \quad (11)$$

Then, $\left\{ W_N^{\beta, \eta} \right\}_N$ is uniformly integrable.

$$\mathbb{E} W_\infty^{\beta, \eta} = \lim_{N \rightarrow \infty} \mathbb{E} W_N^{\beta, \eta} = 1$$

Changing the setup

Our polymer measure

$$\frac{d\mathbf{P}_N^{\beta,\omega}}{d\mathbf{P}_0}(S) = \frac{1}{Z_N^{\beta,\omega}} \left(\prod_{n=1}^N (1 + \beta\omega_{n,S_n}) \right), \quad (12)$$

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Assuming,

$$\begin{aligned} \omega_{0,0} &\geq -1 \quad \mathbb{P}\text{-a.s.}, \\ \mathbb{E}[\omega_{0,0}] &= 0 \text{ and} \\ \mathbb{P}[\omega_{0,0} > x] &\stackrel{x \rightarrow \infty}{\sim} C_{\mathbb{P}} x^{-\gamma}, \text{ for } \gamma \in (1, 2), \end{aligned} \quad (13)$$

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Theorem 1

When the environment satisfies the condition above and if $\gamma \leq \gamma_c$, very strong disorder holds, for all values of $\beta > 0$, in all dimensions $d \geq 1$. In particular, if $d \geq 3$ and $\gamma < \gamma_c$,

$$\lim_{\beta \rightarrow 0} \frac{\log |p(\beta)|}{\log \beta} = \alpha, \quad (14)$$

where $\alpha = \alpha(d, \gamma) := \frac{\gamma(\gamma_c - 1)}{\gamma_c - \gamma}$. Also, if $\gamma = \gamma_c$, we have that,

$$\lim_{\beta \rightarrow 0} \frac{\log |p(\beta)|}{\log \beta} = \infty. \quad (15)$$

Our results

Theorem 2

Assuming the same conditions for the environment and if $\gamma > \gamma_c$, we have that $\beta_c > 0$ for dimensions $d \geq 3$.