

Absorbing-state phase transitions

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XXIII Brazilian School on Probability – July, 2019

Background

Context

Non-equilibrium Statistical Mechanics

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Non-equilibrium Statistical Mechanics

Critical phenomena

- self-similar shapes

- large fluctuations

- long-range correlations

- avalanches

Critical phenomena

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$\beta < \beta_c \rightarrow$ well behaved, short-range correlations

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Critical phenomena

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$\beta = \beta_c \rightarrow$ long range, power laws, intricate behavior

$\beta > \beta_c \rightarrow$ well behaved, short-range correlations

But why do we observe critical behavior outside a controlled environment?

Critical phenomena (cont)

Physics literature:

Late 80's – Self-organized criticality

A system that finds a critical state all by itself

Critical phenomena (cont)

Physics literature:

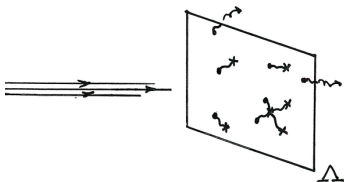
Late 80's – Self-organized criticality

A system that finds a critical state all by itself

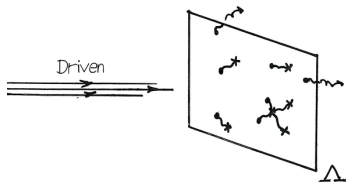
Late 90's – Relate it to ordinary phase transitions

Absorbing-state phase transitions

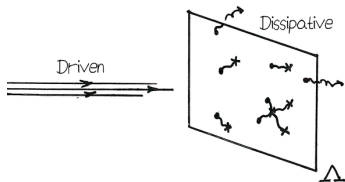
Driven-dissipative dynamics



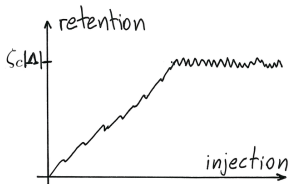
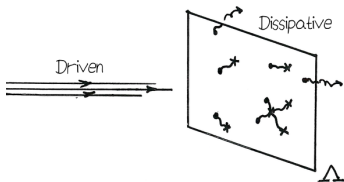
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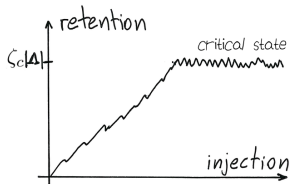
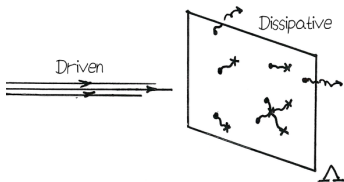
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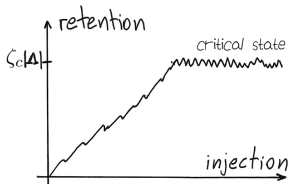
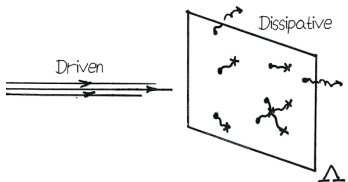
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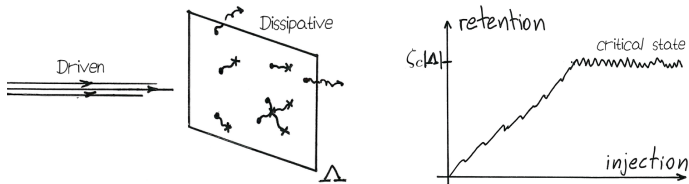


Driven-dissipative dynamics



Infinite-volume conservative system

Driven-dissipative dynamics

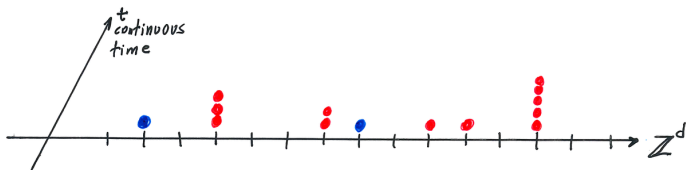


Infinite-volume conservative system

System goes to an absorbing state if the density of particles is below ζ_c , and remains unstable if the density is above ζ_c

Model and predictions

Activated Random Walks



$A \rightarrow S$	rate λ	
$A \text{ jumps}$	1	from x to $x+y$
$AS \rightarrow AA$	∞	with y chosen randomly

Assumptions

- ▶ Jumps to nearest-neighbors

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- ▶ Particles start active
- ▶ At $t = 0$, i.i.d. $\text{Poisson}(\zeta)$ particles
- ▶ $0 < \lambda \leq \infty$

Fixation vs activity

Fixation: each site is eventually stable

Activity: each site is visited infinitely many times

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Dichotomy: either fixation a.s. or activity a.s.

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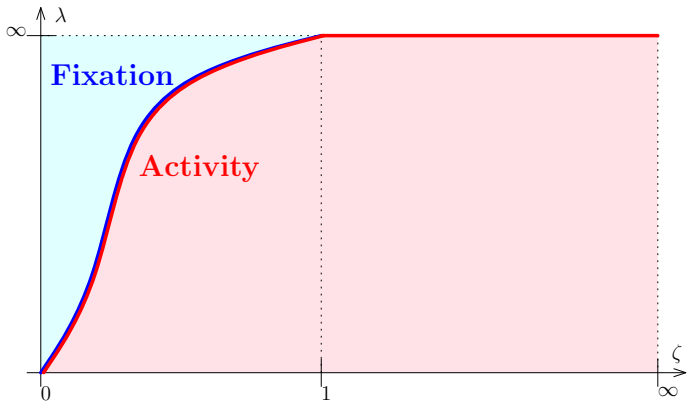
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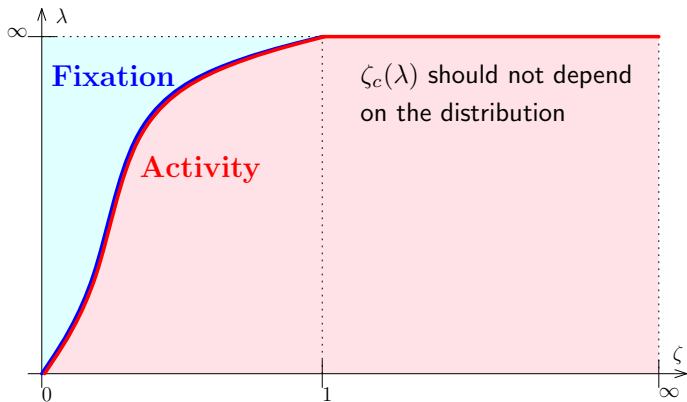
Monotonicity:



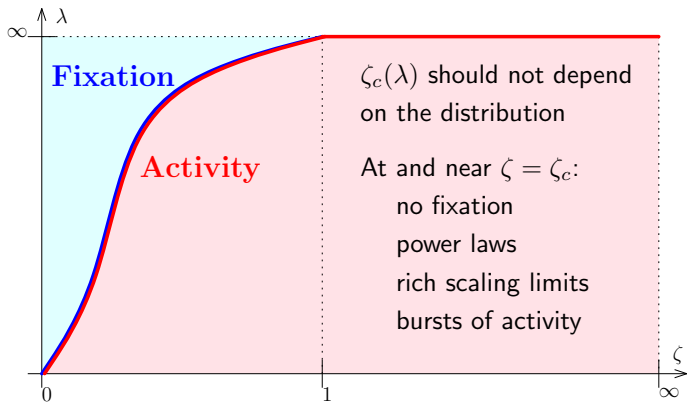
Predictions



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Difficulties

Lack of attractiveness

Overcome by using constructions other than Harris'

Difficulties

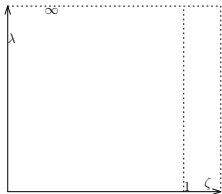
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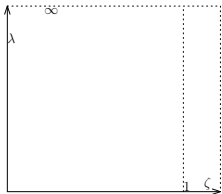
Conservation of particles

Rules out “energy vs. entropy” approaches

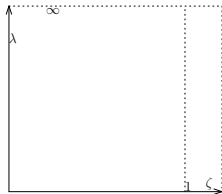
Phase transition results



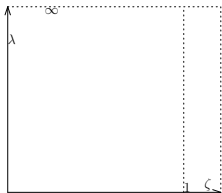
$d = 1$ directed



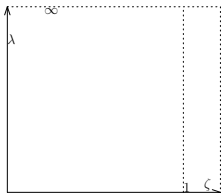
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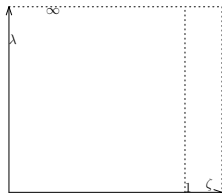
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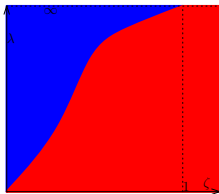
$d \geq 2$ biased



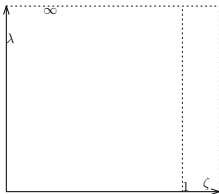
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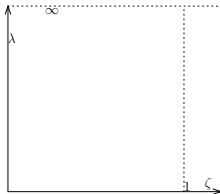
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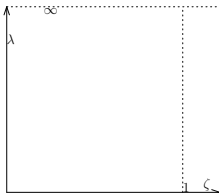
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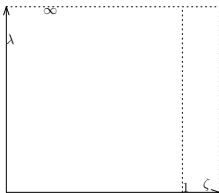
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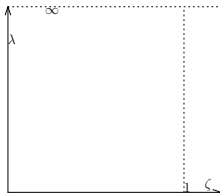
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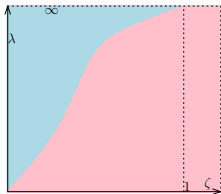


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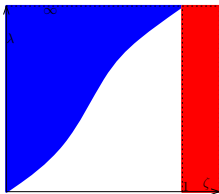


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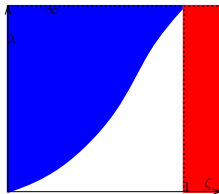
Hoffman, Sidoravicius. *Unpublished* (2004) | Cabezas, R, Sidoravicius. **J Stat Phys** (2014)



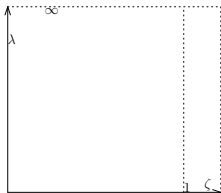
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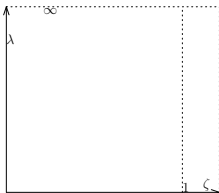
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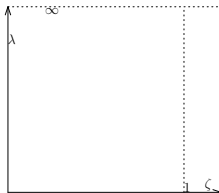
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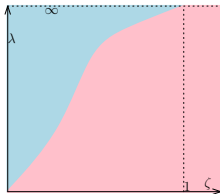
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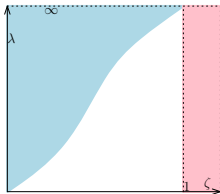
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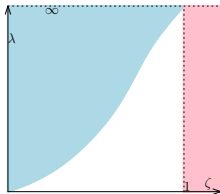
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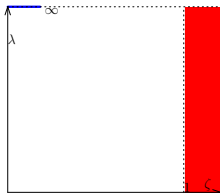
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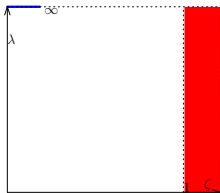
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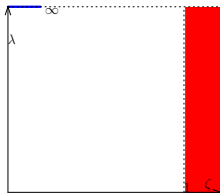
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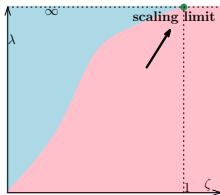
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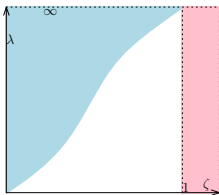
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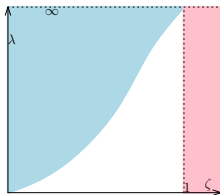
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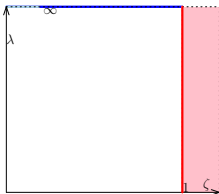
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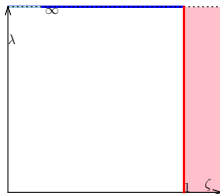
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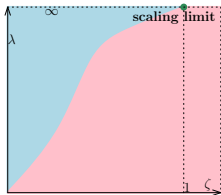
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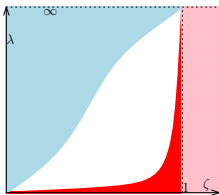
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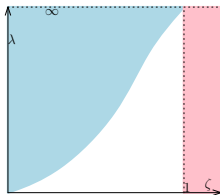
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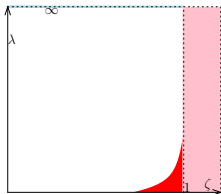
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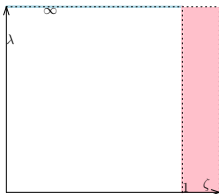
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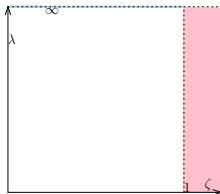
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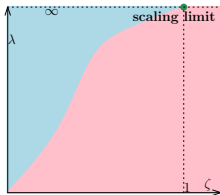
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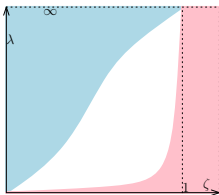
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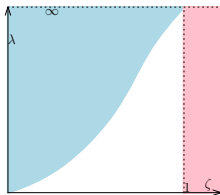
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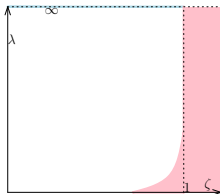
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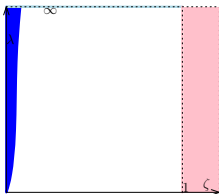
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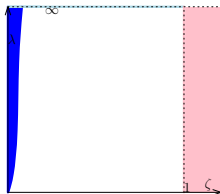
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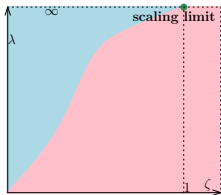
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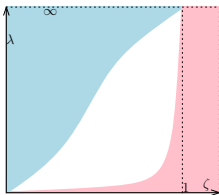
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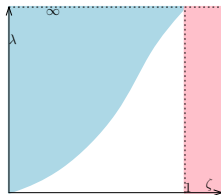
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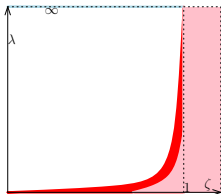
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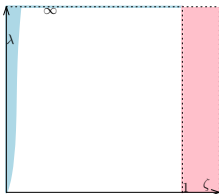
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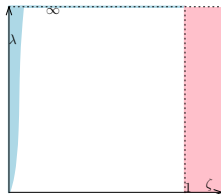
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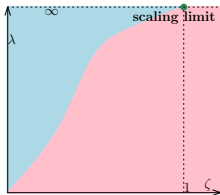
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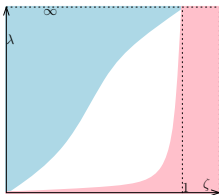
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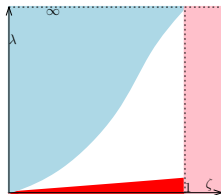
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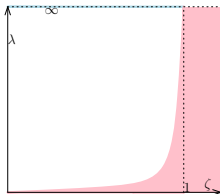
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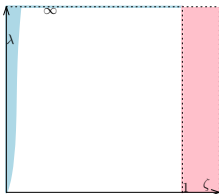
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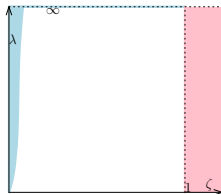
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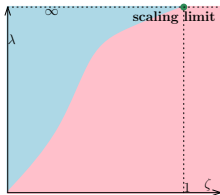
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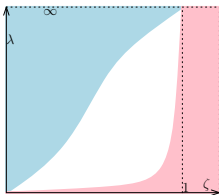
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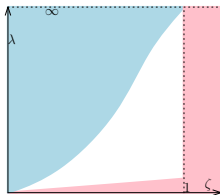
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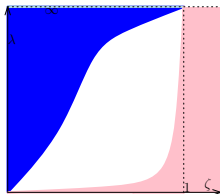
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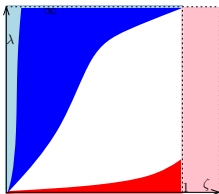
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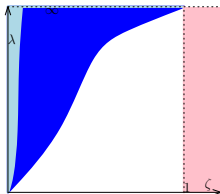
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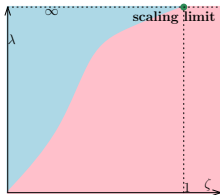
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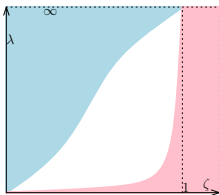
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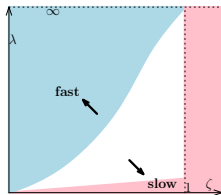
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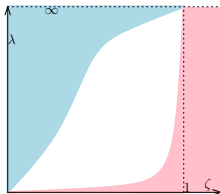
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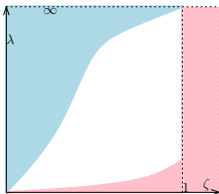
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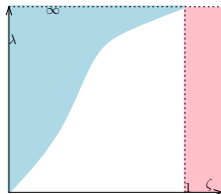
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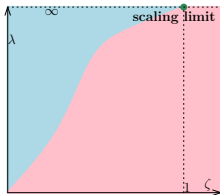
$d \geq 2$ biased



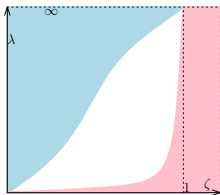
$d \geq 3$ unbiased



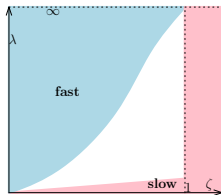
$d = 2$ unbiased



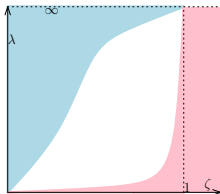
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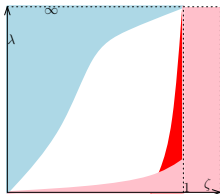
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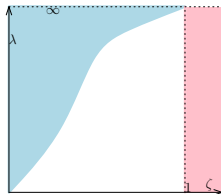
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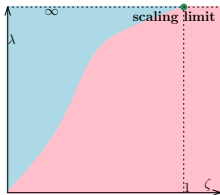


$d \geq 3$ unbiased

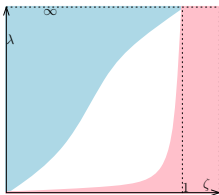


$d = 2$ unbiased

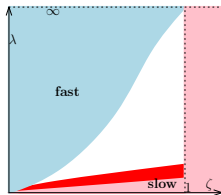
Taggi. *Ann Inst H Poincaré Probab Statist* (2019+)



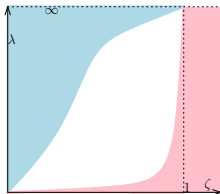
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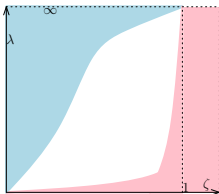
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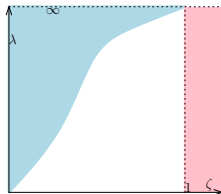
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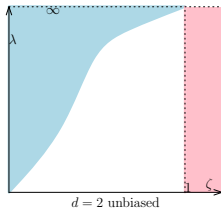
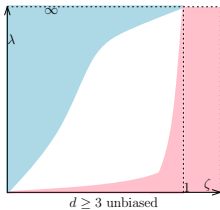
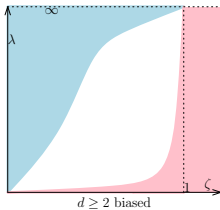
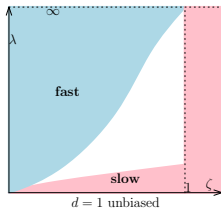
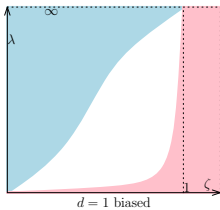
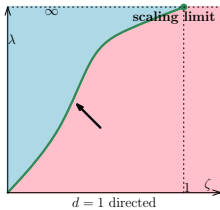


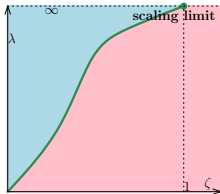
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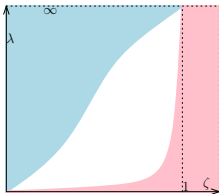
$d = 2$ unbiased

Asselah, R, Schapira. *Writing up*

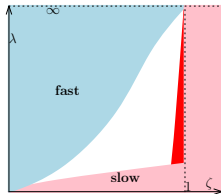




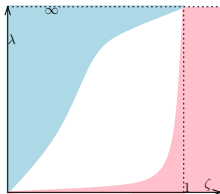
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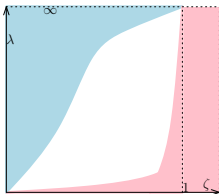
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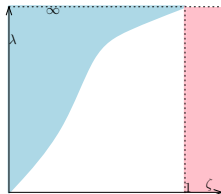
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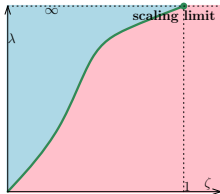
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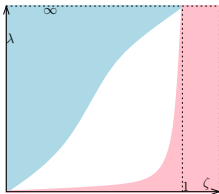
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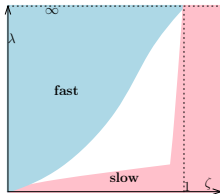
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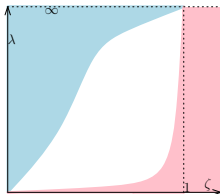
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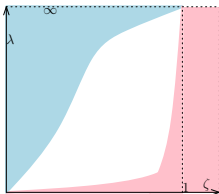
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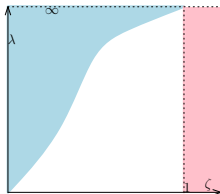
$d = 1$ unbiased



$d \geq 2$ biased



$d \geq 3$ unbiased



$d = 2$ unbiased

Constructions and main tools

Constructions

Harris graphical construction

clocks with marks at each site

Constructions

Harris graphical construction

clocks with marks at each site

Site-wise construction

stack of instructions at each site

Constructions

Harris graphical construction

clocks with marks at each site

Site-wise construction

stack of instructions at each site

Particle-wise constructions

particles start with a life plan and do pause/resume

Main tools

Site-wise representation and Abelianess

Diaconis, Fulton. **Rend Semin Mat Torino** (1991) | Eriksson. **SIAM J Discrete Math** (1996)

Main tools

Site-wise representation and Abelianess

Diaconis, Fulton. **Rend Semin Mat Torino** (1991) | Eriksson. **SIAM J Discrete Math** (1996)

Relate it to the dynamics, preserving monotonicity, etc

Reduce fixation-activity question to *toppling procedures*

R, Sidoravicius. **Invent Math** (2012)

Main tools (cont)

Assuming a particle-wise construction is well-defined:
a particle stays active \Rightarrow sites stay active

Amir, Gurel-Gurevich. **Electron Commun Probab** (2010)

Main tools (cont)

Assuming a particle-wise construction is well-defined:
a particle stays active \Rightarrow sites stay active

Amir, Gurel-Gurevich. **Electron Commun Probab** (2010)

Well-definedness of the particle-wise construction
 \rightarrow ergodicity, mass transport, coupling, surgery

Averaged criterion for activity

R, Tournier. **Ann Inst H Poincaré Probab Statist** (2018)

Criterion for activity

└ Stabilizing a large box forces a large number of particles to **visit a specific site, wpp**

R, Sidoravicius. **Invent Math** (2012)

Criterion for activity

⊢ Stabilizing a large box forces a large number of particles to **visit a specific site, wpp**

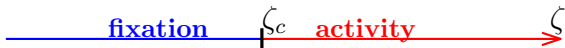
R, Sidoravicius. **Invent Math** (2012)

⊢ Stabilizing a large box forces a positive fraction of the particles to **leave the box, on average**

R, Tournier. **Ann Inst H Poincaré Probab Statist** (2018)

Sharpness

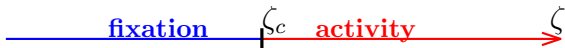
Theorem. Given d, λ, p , there exists ζ_c such that

 ζ_c activity ζ

for all ergodic initial states with density ζ

R, Sidoravicius, Zindy. **Ann Henri Poincaré** (2019)

Theorem. Given d, λ, p , there exists ζ_c such that

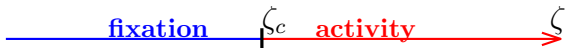
A diagram showing a horizontal line with a vertical tick mark at a point labeled ζ_c . The region to the left of the tick mark is labeled "fixation" in blue text, and the region to the right is labeled "activity" in red text. A red arrow points to the right from the tick mark, ending at a symbol ζ .

for all ergodic initial states with density ζ

R, Sidoravicius, Zindy. **Ann Henri Poincaré** (2019)

– Can drop the previous “i.i.d. Poisson” assumption

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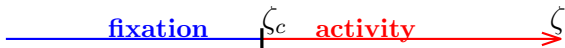


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R, Sidoravicius, Zindy. **Ann Henri Poincaré** (2019)

- Can drop the previous “i.i.d. Poisson” assumption
- Restrictive proofs now yield general theorems
- Contributes to ongoing discussion about some dissipative models mixing better than others

Scaling limits and avalanches

Critical one-dimensional model

$d = 1$, directed walks (integrable case)

i.i.d. initial condition with critical density $\zeta = \zeta_c = \frac{\lambda}{1+\lambda}$

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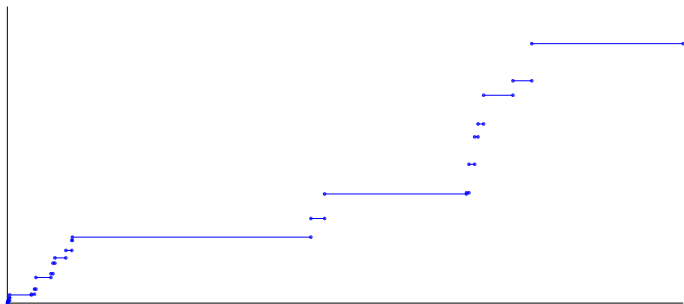
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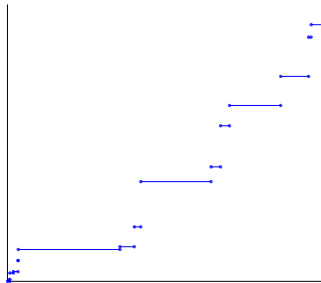
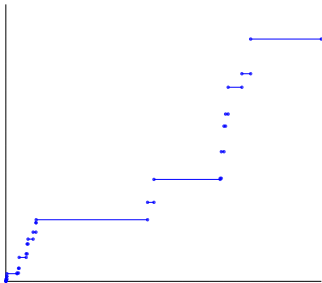
Run the dynamics on $V_n = [0, n]$ until it is stable.

$C_n :=$ how many particles cross the origin.

Released the active particles at $x = n + 1$, let them interact with the leftovers of previous step. $C_{n+1} \geq C_n$

Simulation





Scaling limit

Theorem (Cabezas, R '19).

The counting process $(C_n)_{n \geq 0}$ rescales to $(C_x^\rho)_{x \geq 0}$

Pure-jump process constructed from a collection of correlated reflected coalescing Brownian motions

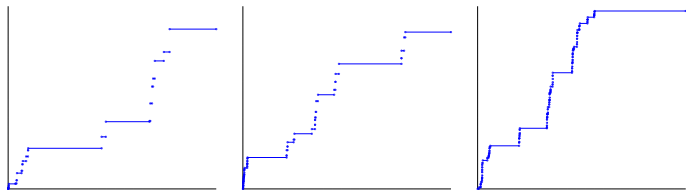
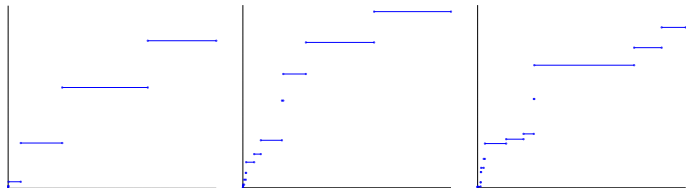
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Correlations depend on $\rho = \frac{\sigma_s}{\sigma_p} \in (0, 1]$, where $\sigma_s^2 = \zeta - \zeta^2$ and σ_p^2 is the variance of initial condition



$\rho = 1.00, 0.50, 0.30, 0.10, 0.05, 0.00$

Notes on some selected proofs

Toppling procedures

Abelian particle-wise construction

$m_V(x) :=$ odometer at x when V is stabilized

$$\lim_k \lim_V \mathbb{P}[m_V(x) \geq k] = \begin{cases} 0 \Leftrightarrow \text{Fixation} & \text{(B)} \\ 1 \Leftrightarrow \text{Activity} & \text{(U)} \end{cases}$$

$M_n :=$ #Particles which quit when B_n is stabilized

$$\limsup_n \frac{\mathbb{E}[M_n]}{|B_n|} > 0 \Rightarrow \text{Activity} \quad \text{(E)}$$

Examples

Thm (Stauffer, Taggi). $\zeta_c \geq \frac{\lambda}{1+\lambda}$

Thm (Taggi; R, Tournier). $\zeta_c \leq F_p(\lambda)$

Particle-wise construction

Labeled particles \rightarrow mass transport, ergodicity, surgery

Construction: assign to each particle a CTRW+beep

MTP Example. Assume particles fixate a.s.

$$\zeta = \mathbb{E}[\text{start at } \mathbf{o}] = \mathbb{E}[\text{settle at } \mathbf{o}] \leq 1$$

Thm. Site fixation $\Rightarrow \zeta < 1$

Particle-wise construction (cont)

Thm. The PWC is well-defined

Add particles one by one, updating the whole evolution

┆ Life of each particle is well-defined through some limit

Main step: $\forall x, T$, the number of particle additions that affect site x by time T has finite expectation

Open problems and questions

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Proofs of fixation/activity that give different insights

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Unbiased walks on \mathbb{Z}^2 : $\zeta_c < 1$ for some $\lambda < \infty$

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Many more...